Topological Kondo Insulator - SmB6

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OUTLINE

MOTIVATION

SmB6 ANALYSIS

• IS IT REALLY TOPOLOGICAL?

MOTIVATION

REMEMBER KONDO INSULATOR BEHAVIOR?



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How about SmB6



How about SmB6



How about SmB6



So Why Are We Interested in SmB6

- Mystery at low temperature (T<3K)
 - => resistivity plateau and anomalous susceptibility behavior

 The strong correlation + strong SO coupling + inversion of the bands with opposite parity
 => interacting TI candidate

• It is a true bulk insulator, unlike the other known TI

SmB6 Analysis - band theory and symmetry

What is SmB6





Ground state config of Sm $[Xe]4f^{6}6s^{2}$

Main role - Sm Supporting role - B

draw the chemical picture

What is SmB6





Conduction electron Sm - 5d + B- 2p Localized electron Sm - 4f

SmB6 Band Configuration - Group Theory



Assuming atomic force > S-O > crystal field

mention the degeneracy and predict the band diagram from here

SmB6 Band Configuration - Group Theory

• Consider the bands near Fermi level.

Choose conduction e : $5d_{J=5/2}\Gamma_8$ localized e : $4f_{J=5/2}\Gamma_8$

 ${\scriptstyle \bullet}$ These states are in the Double group O

 Γ_8 is four fold degeneracy





Model Setup - PAM

 $4f^6:4f^5\approx 3:7$

Model Setup

- ξ_k onsite + n.n. hopping
- ϵ_k onsite + n.n. hopping
- V n.n. hopping anisotropy hybridization gap U_{ff} onsite

• STRATEGY

1. Find the possible Hamiltonian allowed by the symmetry. Ex: rotation or time reversal...

 $H_{sym} + U f^{\dagger} f^{\dagger} f f$

2. Do the slave boson procedure to absorb the repulsion term into effective Hamiltonian

H_{eff}

3. Diagonalize the H and study the band diagram

 E_k

• STEP ONE

Symmetry in SmB6

 $\frac{\pi}{2}$ rotation along x, y, z direction, π reflection to x, y, z plane, TR

Operator transformation

$$\begin{aligned} R^{z}(\pi/2)a_{lms}^{\dagger}R^{z}(\pi/2)^{\dagger} &= e^{-im\pi/2}e^{-is_{z}\pi/4}a_{lms_{z}}^{\dagger} \\ R^{y}(\pi/2)a_{lms}^{\dagger}R^{y}(\pi/2)^{\dagger} &= \Sigma\Lambda_{m'm}e^{-is_{y}\pi/4}a_{lm's_{z}}^{\dagger} \\ R^{x}(\pi/2)a_{lms}^{\dagger}R^{x}(\pi/2)^{\dagger} &= \Sigma e^{im'\pi/2}|\Lambda_{m'm}e^{-im\pi/2}e^{-is_{x}\pi/4}a_{lm's_{z}}^{\dagger} \\ R_{x}a_{lms}^{\dagger}R_{x}^{\dagger} &= \sigma_{x}a_{l,-ms}^{\dagger} \\ R_{y}a_{lms}^{\dagger}R_{y}^{\dagger} &= \sigma_{y}(-1)^{m}a_{l,-ms}^{\dagger} \\ R_{z}a_{lms}^{\dagger}R_{z}^{\dagger} &= \sigma_{z}(-1)^{l+m}a_{lms}^{\dagger} \\ Ta_{lms}^{\dagger}T^{\dagger} &= e^{-i\pi\sigma_{y}/2}(-1)^{m}a_{l-ms}^{\dagger} \end{aligned}$$

• STEP ONE

Hamiltonian respect to Symmetry

1 -

$$OHO^{\dagger} = H \qquad \qquad h(\mathbf{R}) = \begin{pmatrix} h^{d}(\mathbf{R}) & V(\mathbf{R}) \\ V^{\dagger}(\mathbf{R}) & h^{f}(\mathbf{R}) \end{pmatrix},$$

$$\begin{split} H_{d} &= \mu_{d} \sum_{r} d_{r}^{\dagger} d_{r} + t^{d} \sum_{r} d_{r}^{\dagger} \begin{pmatrix} \mathbf{1} & \mathbf{0} \\ \mathbf{0} & \eta_{d} \mathbf{I} \end{pmatrix} d_{r\pm \hat{x}} \\ &+ t^{d} \sum_{r} d_{r}^{\dagger} \begin{pmatrix} (\frac{1}{4} + \frac{3}{4}\eta_{d})\mathbf{I} & (\frac{\sqrt{3}}{4} - \frac{\sqrt{3}}{4}\eta_{d})\mathbf{I} \\ (\frac{\sqrt{3}}{4} - \frac{\sqrt{3}}{4}\eta_{d})\mathbf{I} & (\frac{3}{4} + \frac{1}{4}\eta_{d})\mathbf{I} \end{pmatrix} d_{r\pm \hat{x}} \\ &+ t^{d} \sum_{r} d_{r}^{\dagger} \begin{pmatrix} (\frac{1}{4} + \frac{3}{4}\eta_{d})\mathbf{I} & -(\frac{\sqrt{3}}{4} - \frac{\sqrt{3}}{4}\eta_{d})\mathbf{I} \\ -(\frac{\sqrt{3}}{4} - \frac{\sqrt{3}}{4}\eta_{d})\mathbf{I} & (\frac{3}{4} + \frac{1}{4}\eta_{d})\mathbf{I} \end{pmatrix} d_{r\pm \hat{y}} & H_{f} = H_{d}(d \longrightarrow f) \\ &+ Uf^{\dagger} f^{\dagger} ff f \\ \\ H_{df} &= \sum_{r} d_{r}^{\dagger} \begin{pmatrix} V_{1}\sigma_{z} & \mathbf{0} \\ \mathbf{0} & V_{2}\sigma_{z} \end{pmatrix} f_{r+\hat{z}} - d_{r}^{\dagger}h_{df}(\hat{z})f_{r-\hat{z}} \\ &+ \sum_{r} d_{r}^{\dagger} \begin{pmatrix} -(\frac{1}{4}V_{1} + \frac{3}{4}V_{2})\sigma_{x} & -\frac{\sqrt{3}}{4}(V_{1} - V_{2})\sigma_{x} \\ -\frac{\sqrt{3}}{4}(V_{1} - V_{2})\sigma_{x} & -(\frac{3}{4}V_{1} + \frac{1}{4}V_{2})\sigma_{x} \end{pmatrix} f_{r+\hat{x}} - d_{r}^{\dagger}h_{df}(\hat{x})f_{r-\hat{x}} \\ &+ \sum_{r} d_{r}^{\dagger} \begin{pmatrix} -(\frac{1}{4}V_{1} + \frac{3}{4}V_{2})\sigma_{y} & \frac{\sqrt{3}}{4}(V_{1} - V_{2})\sigma_{y} \\ \frac{\sqrt{3}}{4}(V_{1} - V_{2})\sigma_{y} & -(\frac{3}{4}V_{1} + \frac{1}{4}V_{2})\sigma_{y} \end{pmatrix} f_{r+\hat{y}} - d_{r}^{\dagger}h_{df}(\hat{y})f_{r-\hat{y}} \end{split}$$

• STEP ONE

Hamiltonian in Momentum Space

$$\begin{aligned} H_d &= \mu_d \sum_k d_k^{\dagger} d_k + \frac{t^d}{2} \sum_k d_k^{\dagger} \left(\begin{array}{c} (\phi_1 + \eta_d \phi_2) \mathbf{I} & (1 - \eta_d) \phi_3 \mathbf{I} \\ (1 - \eta_d) \phi_3 \mathbf{I} & (\phi_2 + \eta_d \phi_1) \mathbf{I} \end{array} \right) d_k \\ \text{where } \phi_1 &= \cos k_x + \cos k_y + 4 \cos k_z, \ \phi_2 &= 3(\cos k_x + \cos k_y), \ \phi_3 &= \sqrt{3}(\cos k_x - \cos k_y), \\ H_f &= H_d(d \to f) + U f^{\dagger} f^{\dagger} f f \end{aligned}$$

$$H_{df} = -\frac{i}{2} \sum_{k} d_{k}^{\dagger} \left(\begin{array}{cc} 4V_{1}\tilde{\sigma_{z}} - (V_{1} + 3V_{2})(\tilde{\sigma_{x}} + \tilde{\sigma_{y}}) & \sqrt{3}(V_{2} - V_{1})(\tilde{\sigma_{x}} - \tilde{\sigma_{y}}) \\ \sqrt{3}(V_{2} - V_{1})(\tilde{\sigma_{x}} - \tilde{\sigma_{y}}) & 4V_{2}\tilde{\sigma_{z}} - (3V_{1} + V_{2})(\tilde{\sigma_{x}} + \tilde{\sigma_{y}}) \end{array} \right) f_{k}^{\dagger}$$

where $\tilde{\sigma}_i = \sigma_i \sin k_i$.

• STEP TWO

Absorb the Correlation Term by Slave Boson

$$H_f = U \sum_i \sum_{lpha=1}^4 \sum_{eta
eq lpha} f_{ilpha}^\dagger f_{ilpha} f_{ieta}^\dagger f_{ieta} f_{ieta}$$

consider infinite U, where we can project out all the states with $n_f > 1$, by replacing the bare f field by Hubbard operators X.

$$f_{i\alpha}^{\dagger} \to X_{\alpha 0}(i) = f_{i\alpha}^{\dagger} b_i, \quad f_{i\alpha} \to X_{0\alpha}(i) = b_i^{\dagger} f_{i\alpha},$$

 $\sum_{\alpha=1}^{-} f_{i\alpha}^{\dagger} f_{i\alpha} + b_i^{\dagger} b_i = 1.$

Partition function

$$\begin{split} L &= \sum_{i} b_{i}^{\dagger} \frac{d}{d\tau} b_{i} + \sum_{ij} \sum_{\alpha,\beta=1}^{4} f_{i\alpha}^{\dagger} \left[\delta_{ij} \delta_{\alpha\beta} \left(\frac{d}{d\tau} + \varepsilon_{f} \right) + b_{i} t_{ij,\alpha\beta}^{(f)} b_{j}^{\dagger} \right] f_{j\beta} + \sum_{\mathbf{k}\sigma} \sum_{a,b=1}^{2} c_{a\mathbf{k}\sigma}^{\dagger} \left(\frac{d}{d\tau} + \varepsilon_{ab}^{(d)}(\mathbf{k}) \right) c_{b\mathbf{k}\sigma} \\ &+ \frac{1}{2} \sum_{\langle ij \rangle} \sum_{\mathbf{k}\sigma} \sum_{a=1}^{2} \sum_{\beta=1}^{4} \left(V_{ia\sigma,j\beta} c_{ai\sigma}^{\dagger} b_{i}^{\dagger} f_{j\beta} + \text{h.c.} \right) + \sum_{j} i\lambda_{j} \left(\sum_{\alpha=1}^{4} f_{j\alpha}^{\dagger} f_{j\alpha} + b_{j}^{\dagger} b_{j} - 1 \right) \end{split}$$

After the mean field approximation (saddle point), the min value is at

$$b_{\mathbf{q}}(\tau) = b\delta_{\mathbf{q},0}, \quad i\lambda_{\mathbf{q}}(\tau) = (E_f - \varepsilon_f)\delta_{\mathbf{q},0},$$

Put this back to Lagrangian, and it reduced to $L = L_0 + H_v$

$$\begin{split} L_0(\tau) &= (b^2 - 1)(E_f - \varepsilon_f) + \sum_{\mathbf{k}\alpha} \sum_{n=1}^2 \tilde{f}_{n\mathbf{k}\alpha}^{\dagger} \left(\frac{d}{d\tau} + E_{n\mathbf{k}}^{(f)}\right) \tilde{f}_{n\mathbf{k}\alpha} + \sum_{\mathbf{k}\sigma} \sum_{a=1}^2 d_{a\mathbf{k}\sigma}^{\dagger} \left(\frac{d}{d\tau} + E_{a\mathbf{k}}^{(d)}\right) d_{a\mathbf{k}\sigma} \\ \tilde{H}_V &= i V_{df} \frac{b}{2} \sum_{\mathbf{k}} \hat{d}_{\mathbf{k}\alpha}^{\dagger} [\hat{\Phi}_{\mathbf{k}}]_{\alpha\beta} \tilde{f}_{\mathbf{k}\beta} + \text{h.c.}, \end{split}$$

In short, we will have a 2X2 effective H with renormalized $E_f V$

$$t_f
ightarrow b^2 t_f$$

 $V_{df}
ightarrow bV_{df}.$ $b^2 + \langle n_f
angle = 1$

We can further integrate out the conduction field, and run the saddle point again. Then we can study the Lagrangian with finite temperature.

$$\frac{\partial S_{eff}}{\partial E_f} = 0, \quad \frac{\partial S_{eff}}{\partial b} = 0.$$

$$b^2 - 1 + 2\sum_{i=1}^4 \sum_{\mathbf{k}} \frac{f(\varepsilon_{i\mathbf{k}})n_F(\varepsilon_{i\mathbf{k}})}{\prod_{l \neq i} (\varepsilon_{i\mathbf{k}} - \varepsilon_{l\mathbf{k}})} = 0, \quad n_F(x) = \frac{1}{e^{\beta x} + 1},$$
Mean Field Equation
$$8(E_f - \varepsilon_f) + 2\sum_{i=1}^4 \sum_{\mathbf{k}} \frac{\partial \varepsilon_{i\mathbf{k}}}{\partial E_f} n_F(\varepsilon_{i\mathbf{k}}) = 0.$$

0.12

0.1

• STEP THREE

ENERGY DIAGRAM AND PFT DIAGRAM



The gap is opened by the hybridization of d, f bands

Tuesday, April 22, 14

SmB6 - Is it topological? Really?

Z2 number

• Based on Hasan and Kane's work, a 3D TI insulator is characterize by four \mathbb{Z}_2 topological invariants $(\nu_0; \nu_1 \nu_2 \nu_3)$

strong TI index
$$I_{STI} = (-1)^{\nu_0} = \prod_{m=1}^8 \delta_m = \pm 1$$
weak TI index $I^a_{WTI} = (-1)^{\nu_a} = \prod_{k_m \in P_j} \delta_m = \pm 1$ $\delta_m = \operatorname{sgn}(\xi_{\mathbf{k}^*_m} - \tilde{\varepsilon}^{(f)}_{\mathbf{k}^*_m})$ $\mathbf{k}^*_{\mathbf{k}^*_m}$ Example
STI (1;0,0,0)
WTI (0;1,1,1) $\mathbf{k}^*_{\mathbf{k}^*_m}$

Z2 number for SmB6



Z2 number = (1; 1, 1, 1)

Z2 number for SmB6



 $v = 3 - n_f$

- v gives the valence of the Sm ion, while n_f measures the number of fholes in the filled $4f^6$ state, so that $n_f = 1$ corresponds to the $4f^5$ configuration.
- SmB6 is mixed valence materials. $4f^{5.7}$ v = 3 - 0.3 = 2.7

How does the experiment say?

Transport

ARPES



How does the experiment say?

• Other SS?



How does the experiment say?

• Some other possible SS... polarity.. band bending Fermi gas.. more



More experiments like MR, QO, and magnetic doping Some support topo.
 SS, and some against it...

Review and Question

- The Kondo insulator which is a true insulator has strong electron correlation and spin-orbital coupling => An ideal TI candidate
- The PAM + symmetry constraint + slave boson => effective H
- The energy diagram and the PFT result show the existence of gap and edge state
- The Z2 number = -1. Theoretical prediction of TKI. However, experiments are needed to verify it.

Question

- Can Kondo insulator adiabatically connect to normal insulator?
- Does U help to increase the TI phase?
- How does SmB6 behave under different T, P, B ? And the phase diagram?
- The Kondo screening, spin configuration in SmB6?
- Is the edge state really topological? What is the origin of the S.S.

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