Spin-helical transport in normal and superconducting topological insulators

Introduction

- Spin-momentum locking=helical edge states
- Helical edge states lead to novel transport properties
- But its not as easy to detect as for Quantum Spin Hall
- Requires Fermi energy in the gap





- Weak antilocalization
- Fractional Quantum Hall
- Interface with Superconductors

- Hamiltonian: $\mathcal{A}\sigma \cdot k$ term (phenomenological)
- very strong spin-orbit coupling leads to $E(k) \approx \mathcal{A}k$
- Impurity protection=diffusive regime=constant spin
- Q: How does large A protect against impurities and enforce time-reversal symmetry?

3.2 Weak antilocalization

- In diffusion regime, a lot like classical (Drude) insulator
- Spin-orbit coupling leads to destructive interference when backscattered
- Increases conductivity
- Magnetic Field can change the conductivity in well-defined ways
- Due to small penetration of edge state into the bulk, field perpendicular to the surface has a much larger effect than field parallel to the surface



- Take only spin-orbit term of Hamiltonian, add magnetic field
- Q: Why doesn't the strong magnetic field destroy the topological order?
- Can use gauge freedom to set diagonal terms to 0, off-diagonal to $-i\hbar k_x eBx = a$
- Diagonalize
- Kubo formula: $\langle j_{\alpha}(t) \rangle = \int dt i \theta(t-t') \langle [j_{\alpha}(t), j_{\beta}(t')] \rangle A(t')$
- Fourier transform!
- $\bullet\,$ Note that j's are in the interaction picture, and $E=A/\omega$

$$\begin{array}{l} j_{\alpha} = e v \sigma_{\alpha} \\ \langle \bar{n'} | \sigma_{\pm} | \bar{n} \rangle = \pm i \frac{1}{2} \delta_{|n'|,|n| \pm 1} \end{array}$$

•
$$\sigma_{xy}(\omega) = \sum_{n} \sum_{n'} (f_n - f_{n'}) \frac{|\langle \bar{n'} | \sigma_+ | \bar{n} \rangle|^2 - |\langle \bar{n'} | \sigma_- | \bar{n} \rangle|^2}{(\epsilon_n + \epsilon_{n'})^2 / \hbar^2 + (\tau - i\omega)^2}$$

- Terms like $-N 1 \rightarrow N + 2$ and $-N 2 \rightarrow N + 1$ cancel
- So for equal number of particles and holes, only $N \rightarrow N + 1$ and $-N \rightarrow N + 1$ contribute

• Leads to
$$\sigma_{xy} = \frac{e^2}{\hbar} \left[N + \frac{1}{2} \right]$$

• For different number of particles and holes, can have another contribution which breaks up half-integer effect





- Top and bottom layer can have different *N*, makes it difficult to observe this effect
- Can also compute AC conductivities

- \bullet Novel behaviour expected at S/TI/S junction
- Spin-momentum locking induces antisymmetric p-wave pairing on the superconductor
- Could form Majoranas at interface, but for thin TI these might be able to pair
- Direct calculation confirms existence of p states
- Andreev Bound States (ABS) are states in the junction formed by Andreev reflection
- When the states are perpendicular to the boundary, there is no backscattering
- This leads to different states (and different currents through the junction) than in the non-topological case